

Institute of Theoretical Physics
University of Wroclaw
Wroclaw, Cybulskiego 36

Wroclaw, March 1983
Preprint No 577

CLASSICAL LIMIT OF THE CANONICAL ANTICOMMUTATION RELATIONS
AND SELF DUALITY OF THE INFINITE DIMENSIONAL GRASSMANN ALGEBRA^(*)

by

A.Z. Jadczyk & K. Pilch

Abstract A classical limit of the twisted group algebra for canonical anticommutation relations (CAR) is constructed and shown to exhibit a kind of selfduality relevant for applications in supersymmetric models.

(*) To appear in "Selected Topics on Quantum Theory of Fields and Particles: A Birthday Volume to Jan Lopuszanski" Ed. by B. Jancewicz and J. Lukierski, World Sci. Publ. Co. Singapore.

Symb. 18. (2.3)

1. Introduction

When dealing with Fermi systems it is often convenient to introduce an auxiliary Grassmann algebra Q . Originally Q was introduced to accomodate "eigenvalues" of anticommuting operators of CAR so that a formal analogy between Feynmann integrals for Bose and Fermi systems could be established (see [1,2]). Later on the canonical formalism of classical mechanics was further extended by adding Grassmann coordinates in the configuration space [2-6]. With the invention of supersymmetry [7-9] and establishing its quantum theoretical basis [10] an auxiliary Grassmann algebra was widely applied for integrating of Lie superalgebra relations. Although it is true that one can deal with a quantum theory of supersymmetric systems without "anticommuting variables" (see e.g.[10,11]), nevertheless a possibility that Q can be an indispensable mathematical tool (like the imaginary numbers) should be seriously taken into account. So far, however, the mathematics of Q -numbers is in its early stage of development. In [12] differential supergeometry was developed on a basis of a self-dual Banach-Grassmann algebra Q . In [13] an infinite dimensional algebra (called B_∞ there) was defined. In the present paper we define Q as a classical limit of the CAR twisted group algebra and show that Q defined in such a way is selfdual according to the requirements of [12], and coincides with the B_∞ -algebra of A. Rogers [13].

2. GRASSMANN ALGEBRA AS A CLASSICAL LIMIT OF CAR

Let N denotes the set $\{1, 2, \dots, N\}$ or \mathbb{N} for $N = \infty$. Let

$$a_i a_j + a_j a_i^* = 2\delta_{ij}, \quad a_i = a_i^*, \quad i, j \in N \quad (2.1)$$

be a representation of the Canonical Anticommutation Relations (CAR) with $\frac{N}{2} < \infty$ degrees of freedom. It is natural to associate with (2.1) the group J consisting of all finite subsets $I = \{i_1, \dots, i_k\}$,

$i_1 < i_2 < \dots < i_k$ of the set N with the symmetric difference Δ as the group law. The unitary operators

$$U(I) := a_{i_1} \dots a_{i_k}, \quad I = \{i_1, \dots, i_k\} \in J, \quad (2.2)$$

form a projective representation of the group J ([14,Ch.XIV,§ 8], see also [15, Sect. 3.10])

$$U(I)U(I') = \omega(I, I')U(I \Delta I') \quad (2.3)$$

with

$$\omega(I, I') = (-1)^{P(I, I')}, \quad P(I, I') = \sum_{j \in I'} P(I, j), \quad (2.4)$$

$$P(I, j) = (\text{the number of } i \in I \text{ such that } i < j). \quad (2.5)$$

The multiplier ω satisfies the standard cocycle relations

$$\omega(e, I) = \omega(I, e) = 1, \quad (2.6)$$

$$\omega(I, I')\omega(I \Delta I', I'') = \omega(I', I'')\omega(I, I' \Delta I''), \quad (2.7)$$

where $e = \emptyset$ is the identity element of the group J . We can also define a one parameter family of multipliers

$$\omega_\lambda(I, I') = \lambda^{|I \cap I'|} \omega(I, I'), \quad \lambda \in [0, 1], \quad (2.8)$$

where

$$|I \cap I'| = (\text{the number of elements in } I \cap I').$$

The ω_λ 's may be considered as ω -multipliers (see (2.2) and (2.3))

for the modified CAR

$$a_i(\lambda)a_j(\lambda) + a_j(\lambda)a_i(\lambda) = 2\lambda\delta_{ij}. \quad (2.9)$$

In particular the cocycle relations (2.6), (2.7) are satisfied for all λ and

$$|\omega_\lambda(I, I')| \leq 1, \quad \lambda \in [0, 1]. \quad (2.9a)$$

Using ω_λ one can construct for J the so called twisted group algebra [16] which is a generalization of the ordinary group algebra. Let μ be the invariant measure on J normalized to $\mu(I) = 1$, $I \in J$. The twisted group algebra $Q^N(\lambda)$ is defined as $L^1(J, \mu)$ with the multiplication

$$(f * g)(K) = \sum_{I \in J} f(I)g(I \Delta K)\omega_\lambda(I, I \Delta K), \quad (2.10)$$

$$f, g \in L^1(J, \mu).$$

Since ω_λ are real, Q^N can be real or complex. From now on we assume it is real. Using (2.6), (2.7) and (2.9a) one proves that Q^N is a Banach algebra in the $L^1(J, \mu)$ -norm

$$\|f\| = \sum_{I \in J} |f(I)|, \quad f \in Q^N(\lambda). \quad (2.11)$$

Definition 2.1. The algebra $Q^N = Q^N(\lambda=0)$ is the classical limit for CAR.

It is important to notice that the limit $\lambda \rightarrow 0$ can not be taken on representing operators $a_i(\lambda)$, the limiting procedure is performed on the algebra structure of $L^1(J, \mu)$.

The following Proposition gives a full characterization of Q^N .

Proposition 2.1. The algebra Q^N is uniquely characterized by the following properties:

There exists a set of generators $e_i \in Q^N$, $i \in N$, such that

- (i) $e_i e_j + e_j e_i = 0$, $i, j \in N$,
- (ii) every $f \in Q^N$ is uniquely represented by an absolutely convergent series

$$f = \sum_{I \in J} f^I e_I, \quad \sum_I |f^I| < \infty, \quad f^I \in \mathbb{R},$$

where J is the set of all finite subsets $I = \{i_1, \dots, i_k\}$, $i_1 < \dots < i_k$ of N , and

$$e_I = e_{i_1} \dots e_{i_k}, \quad e_\emptyset = 1,$$

- (iii) for $f \in Q^N$

$$\|f\| = \sum_{I \in J} |f^I|.$$

Remark. In this form the algebra Q^N was introduced by A. Rogers [13].

Proof. Let us denote $\omega_{\lambda=0}(I, I') = x(I, I')$. Then (2.8) implies

$$x(I, I') = \begin{cases} 0 & \text{for } I \cap I' \neq \emptyset, \\ (-1)^P(I, I') & \text{for } I \cap I' = \emptyset, \end{cases} \quad (2.12)$$

and, by (2.4), (2.5), in particular

$$x_{\{i\},\{j\}} = \begin{cases} +1 & \text{for } i < j, \\ 0 & \text{for } i = j, \\ -1 & \text{for } i > j. \end{cases} \quad (2.13)$$

Let e_I be the element of Q^N defined by

$$e_I^{(J)} = \begin{cases} 0 & \text{for } J \neq I, \\ 1 & \text{for } J = I. \end{cases} \quad (2.14)$$

It follows from (2.11) that

$$\|e_I\| = 1, \quad (2.15)$$

and (2.12) gives

$$e_I e_J = e_I^* e_J = \begin{cases} 0 & \text{if } I \cap J \neq \emptyset, \\ x_{\{I,J\}} e_{I \cup J} & \text{if } I \cap J = \emptyset. \end{cases} \quad (2.16)$$

This implies that for $e_i = e_{\{i\}}$

$$e_i e_j = -e_j e_i, \quad (2.17)$$

and for $I = \{i_1, \dots, i_k\}$, $i_1 < \dots < i_k$, we have

$$e_I = e_{i_1} \dots e_{i_k}. \quad (2.18)$$

The proposition follows from an observation that every $f \in Q^N$ can be uniquely written as

$$f = \sum_{J \in \mathcal{J}} f(J) e_J, \quad (2.19)$$

the series being convergent in $L^1(J, \mu)$.

The algebra Q^N admits a natural \mathbb{Z}_2 -grading $Q^N = Q_0^N \oplus Q_1^N$, where Q_1^N is the closed linear subspace generated by odd products of generators e_i , and $Q_0^N = \mathbb{R} \oplus Q_0^N$ is the direct sum of the number field \mathbb{R} and the closed subalgebra Q_0^N generated by even products of elements of Q_1^N . The following norm decompositions hold

$$\|a_0 + a_1\| = \|a_0\| + \|a_1\|, \quad a_r \in Q_r^N, \quad r=0,1, \quad (2.20)$$

$$\|\lambda + a'_0\| = |\lambda| + \|a'_0\|, \quad \lambda \in \mathbb{R}, \quad a'_0 \in Q_0^N. \quad (2.21)$$

The real part of an element $a \in Q^N$ will be denoted by $\sigma(a)$. Therefore $\sigma(\lambda + a') = \lambda$ for $\lambda \in \mathbb{R}$ and $a' \in Q_0^N = Q_0' \oplus Q_1'$.

3. SELF DUALITY OF THE INFINITE DIMENSIONAL GRASSMANN ALGEBRA Q^∞

When N is finite Q^N is 2^N -dimensional and coincides with the standard Grassmann algebra of \mathbb{R}^N . The Proposition 2.1 implies then

Proposition 3.1. If $N < \infty$ then

(i) There exists $e \in Q^N$ such that for every $a \in Q_1^N$

$$ae = ea = 0$$

(ii) For every $a_1, \dots, a_{N+1} \in Q_1^N$

$$a_1 a_2 \dots a_{N+1} = 0$$

The fact that for finite N the equation $ax = 0$, $a \in Q_1^N$, admits non-zero solutions makes it impossible to define partial derivatives

of supersmooth function (see [12]) $f: Q_1^N \rightarrow Q^N$ in such a way that the Leibnitz rule is satisfied (even modulo an ideal). The "superanalysis" formulated in [12] is based on the assumption of selfduality of Q . We shall now prove that the algebra $Q = Q^\infty(\lambda=0)$ defined in Section 2 is selfdual in the required sense.

First of all we observe that by the properties (ii) and (iii) of the Proposition 2.1. for every $a \in Q$ we have

$$\|a\| = \lim_{i \rightarrow \infty} \|e_i a\| = \lim_{i \rightarrow \infty} \|ae_i\|. \quad (3.1)$$

It follows that for each $x \in Q$, if $ax = 0$ (resp. $xa = 0$) for all $a \in Q_1$, then $x = 0$.

Let Q_r^* denote the space of continuous Q_0 -linear maps from Q_r ($r=0,1$) to Q . Every element $q \in Q$ determines $F_q \in Q_r^*$ by

$$F_q(a) = aq, \quad a \in Q_r. \quad (3.2)$$

Since Q_0 is an algebra with unit, it follows that every $F \in Q_0^*$ is of the form F_q for $q = F(1)$.

Proposition 3.2. (Selfduality of Q). For each continuous Q_0 -linear map $F: Q_1 \rightarrow Q$ there exists a unique element $q \in Q$ such that

$$F(a) = aq$$

for all $a \in Q_1$.

To prove the above property of Q we first establish two Lemmas.

Lemma 3.1. Let for each $a \in Q$, $J(a) = \{I \in J: a^I \neq 0\}$, and let $I_0 \subset a$ (resp. $I_0 \perp a$) denotes the fact that $I_0 \not\subset I$ (resp. $I_0 \cap I = \emptyset$) for all $I \in J(a)$. Then

a) $I_0 \not\subset a$ implies

$$\|ae_{I_1} + be_{I_0}\| = \|ae_{I_1}\| + \|be_{I_0}\|$$

for $I_1 \cap I_0 = \emptyset$, and for all $b \in Q$. If $I_0 \perp a$ then $I_0 \not\subset a$ and $\|ae_{I_0}\| = \|a\|$.

b) if $a \in Q$, $I \in J$ and $ae_i = 0$ for all $i \in I$, then there exists a unique $v \in Q$ with $I \perp v$ and such that $a = ve_I$. We also have $\|a\| = \|v\|$.

Proof Follows in a straightforward way from the Proposition 2.1. ■

Lemma 3.2. There is a one-to-one correspondence between continuous Q_0 -linear maps $F: Q_1 \rightarrow Q$ and bounded sequences (f_i) from Q satisfying

$$f_i e_j + f_j e_i = 0 \quad \text{for all } i, j \in \mathbb{N} \quad (3.3)$$

Moreover, $\|F\| = \sup_i \{ \|f_i\| : i \in \mathbb{N} \}$ and $\sigma(f_i) = 0$.

Proof If $F: Q_1 \rightarrow Q$ is Q_0 -linear then, with $x, y, z \in Q_1$, we have $F(x)yz = F(xyz) = -F(yxz) = -F(y)zx$, and therefore $F(x)y = -F(y)x$ for all $x, y \in Q_1$. With $x = e_i$, $y = e_j$ and $f_i = F(e_i)$ we get (3.3). In particular $f_i e_i = 0$ so that $\sigma(f_i) = 0$. Conversely given a bounded sequence $f_i \in Q$ satisfying (3.3) we define $F(e_I) = f_i e_{i_1} \dots e_{i_k}$ and $F(a) = \sum_I a^I F(e_I) : I \in J \}$. It is immediate to see that F so defined is Q_0 -linear and continuous. Since

$$\|F(a)\| = \lim_i \|F(a)e_i\| = \lim_i \|F(e_i)a\| \leq \sup_i \{ \|f_i\| \cdot \|a\| : i \in \mathbb{N} \}$$

we also get $\|F\| = \sup_i \{ \|f_i\| : i \in \mathbb{N} \}$. ■

Proof of the Proposition 3.2. Let $F: Q_1 + Q$ be Q_0 -linear and continuous. Then, according to Lemma 3.2, the sequence $f_i \in Q^*$; $f_i = F(e_i)$ is bounded and $f_i e_j + f_j e_i = 0$. To prove the Proposition we must show that there exists $q \in Q$ such that $f_i = q e_i$ for all $i \in \mathbb{N}$. We shall define q as the limit of a sequence of its approximations q_n constructed as follows:

i) since $f_1 e_1 = 0$ it follows (Lemma (3.1.b)) that

$$f_i = q_1 e_1 \quad \text{with } \{i\} \not\subseteq q_1$$

ii) suppose we have already constructed q_n such that $f_i = q_n e_i$ for $i \in \{1, \dots, n\}$ and $\{1, \dots, n\} \not\subseteq q_n$. Then, since

$$f_{n+1} e_i = -f_i e_{n+1} = q_n e_{n+1} e_i \quad \text{and} \quad f_{n+1} e_{n+1} = 0, \text{ we have}$$

$$(f_{n+1} - q_n e_{n+1}) e_i = 0 \quad \text{for } i \in \{1, \dots, n+1\}.$$

Therefore (Lemma 3.1.b) there exists v_n with $\{1, \dots, n+1\} \setminus v_n$ and such that $f_{n+1} - q_n e_{n+1} = v_n e_1, \dots, e_{n+1}$. We define now $q_{n+1} = q_n + v_n e_1, \dots, e_n$. It is immediate that $q_{n+1} e_i = f_i$ for $i \in \{1, \dots, n+1\}$ and, since $\{1, \dots, n\} \not\subseteq q_n$, we also have $\{1, \dots, n+1\} \not\subseteq q_{n+1}$.

Now, by Lemma 3.1a,

$$\|f_{n+k}\| = \|q_{n+k} e_{n+k}\| = \|q_{n+k-1} e_{n+k} + v_{n+k-1} e_1 \dots e_{n+k}\| =$$

$$+ \|q_{n+k} e_{n+k}\| + \|v_{n+k-1} e_1 \dots e_{n+k}\| >$$

$$> \|q_{n+k-1} e_{n+k}\| > \dots > \|q_n e_{n+k}\|.$$

Therefore $\|q_n\| < \sup\{\|f_{n+k}\| : k \in \mathbb{N}\} < \|f\|$. But $\|q_{n+1}\| = \|q_n + v_n e_1 \dots e_n\| = \|q_n\| + \|v_n\| = \dots = \|q_1\| + \sum_{k=1}^n \|v_k\|$, so that the series $\sum \|v_n\|$ is convergent. Since $\|q_{n+k} - q_n\| = \|\sum_{l=n}^{n+k-1} v_l\| < \sum_{l=n}^{n+k-1} \|v_l\|$, it follows that $q = \lim q_n$ exists in Q . It is also evident that $f_i = q e_i$ for all $i \in \mathbb{N}$. ■

References

- [1] P.T. Matthews, A. Salam, Nuovo Cimento, 2(1955)120.
- [2] J.L. Martin, Proc.Roy.Soc.251A(1959)536.
- [3] F.A. Berezin and M.S. Marinov, JETP Letters 21(1975)321.
- [4] R. Casalbuoni, Nuovo Cimento 33A(1976)389.
- [5] A. Barducci, R. Casalbuoni and L. Lusanna, Nuovo Cimento 35A(1976)377.
- [6] F.A. Berezin and M.S. Marinov, Ann.Phys.104(1977)336.
- [7] Yu. A. Gol'fand, E.P. Likhtman, JETP Letters 13(1971)323
- [8] D.V. Volkov, V.A. Soroka, Phys.Lett.46B(1973)109.
- [9] J. Wess, B. Zumino, Nucl.Phys.B70(1974)39.
- [10] R. Haag, J. Łopuszański, J. Sohnius, Nucl.Phys.B88(1975)255.
- [11] H. Nicolai, "Supersymmetry without anticommuting variables" CERN Preprint TH 2848/80.
- [12] A. Jadczyk, K.Pilch, Commun.Math.Phys.78(1981)373.
- [13] A. Rogers, J.Math.Phys.21(1980)1352.
- [14] S. Lang, "Algebra", Addison-Wesley Publishing Company, Reading, Mass., 1965.
- [15] J. Slawny, Commun.Math.Phys.24(1972)151.
- [16] C.M. Edwards, J.T. Lewis, Commun.Math.Phys.13(1969)119.